Evolution of Galaxies: Review of Stellar Evolution

Observatoire astronomique de Strasbourg

J.Köppen joachim.koppen@astro.unistra.fr

<http://astro.u-strasbg.fr/~koppen/JKHome.html>

Star = self-gravitating ball of gas powered by thermonuclear fusion

Star formation

- Complicated, poorly understood
- Happens in dense, dusty, cool, molecular clouds
- Basic mechanism: Jeans-instability: a selfgravitating gas sphere will collapse if

 E_{thermal} \langle E_{grav} $M/m kT < G M^2/R$ $M \rightarrow M$ jeans $=$ (kT Gm $3^{3/2}$ ($4\pi\rho$ 3 $)^{-1/2}$ $= 1000$ Msun (\overline{T} 10K $3^{3/2}$ (ρ 10^{-24} g/cm3 $)^{-1/2}$

Star formation

- Time scale for initial collapse: free-fall time scale: $\tau_{\text{ff}} \sim 1/\sqrt{G \rho}$
- → Denser clumps will collapse faster than the rest \rightarrow fragmentation

- Collapse is
	- Resisted by magn.field, turbulence, rotation
	- Assisted by external pressure (SN shock, spiral arm)

Basic problem of star formation

How to get rid of angular momentum?

Helpful: -- formation of planet system (?!) -- stellar winds, bipolar molecular outflows, jets

Time scales (Sun)

- Free fall: 1 hr
- Sound wave crossing time = response to pressure imbalance: $R/c_s \sim 1$ hr
- Kelvin-Helmholtz = lifetime of a collapsing star: $E_g/L = GM^2/RL \sim 10^7$ yrs shorter than age of Earth
	- gravitational contraction is NOT main energy source
	- if no other energy sources, star evolves with this timescale
- Nuclear burning: ~10¹⁰ yrs freefall = sound $\lt\lt$ KH $\lt\lt$ nuclear
- **→ Stars evolve in hydrostatic equilibrium**

Static stellar structure

- Assumptions:
	- Single star
	- No magnetic fields
	- No rotation
	- No stellar winds
- Conservation of Mass, Momentum, Energy

Gravity only

Spherically symmetric (1D)

Conservation of Mass

Define mass coordinate: mass within sphere of radius r

$$
M_r = \frac{4\pi}{3} \int_0^r \rho(r') dr'
$$

{ conservation of nuclear species (abundance eⁱ by mass) and with reactions $A+B \rightarrow C$ and rate coefficients $R_{AB,C}$

$$
\varepsilon_i \sum_k \sum_l \varepsilon_k R_{ik,l} = \sum_k \sum_l \varepsilon_k \varepsilon_l R_{kl,i}
$$

consumption = synthesis

Conservation of momentum

Hydrostatic equilibrium:

$$
\frac{dp}{dr}=-\frac{GM_r\rho}{r^2}
$$

 $Pressure p = p$ thermal + P radiative and $\approx aT^{4}/3$

Conservation of Energy

• Energy production

 dL_r \overline{dr} $= 4\pi r^2 \rho e$ Luminosity of sphere r

Energy production rate

- Energy transport
	- $-$ By radiation (interior is highly opaque \rightarrow diffusion approximation):

$$
\frac{dT}{dr} = -\frac{3}{16\pi\sigma c} \frac{K\rho}{T^3} \frac{L_r}{4\pi r^2}
$$

Energy transport: convection

• occurs when layers become dynamically unstable (Schwarzschild criterion for

chem. homog.): if

\n
$$
\frac{\partial \ln T}{\partial \ln p} \Big|_{\text{rad}} < \frac{\partial \ln T}{\partial \ln p} \Big|_{\text{adiab.}} \frac{\gamma - 1}{\gamma}
$$
\n
$$
\frac{dT}{dr} = -\frac{\gamma - 1}{\gamma} \frac{T}{p} \frac{dp}{dr} \qquad \text{with} \quad \gamma = \frac{c_P}{c_V}
$$

- Convection is linked with transport of matter in the unstable regions: Mixing; hot uprising material can overshoot $Ti > T$
	- Physically consistent description NOT available (use recipes obtained from numerical simulations …

Convection: simulation

IcOgh: $time=4300 s$ $v_{\text{Arms,max}}=16.2 km/s$

Stellar Structure equations

- Properties of the material
	- Equation of state (EOS) $p(\rho,T,\varepsilon)$
	- Opacity (LosAlamos, OPAL, ..) $\kappa(\rho,T,\varepsilon)$
	- Energy production $e(\rho,T,\epsilon)$
- Boundary conditions
	- $-$ Centre: Mr(0) = 0, Lr(0) = 0
	- Surface: $T(R) \approx 0$, $p(R) \approx 0$ (= stellar atm.)
- All these equation fully describe the internal structure of a star
- Stellar evolution: sequence of static models with different composition due to nuclear processes

Opacity 'mountain'

Opacity 'mountain'

Solar atmosphere model

TABLE 4.—There are 140 solar abundance models, 72 1/10 solar abundance models, and 72 1/100 solar abundance models listed two per page. The heading for each model gives the effective temperature, log surface gravity, log metal abundance relative to solar, and whether the model is convective plus radiative or purely radiative. The 40 depths actually used in the computation are listed, roughly equally spaced in log τ_{Ross} from -4.5 to 2.0 in steps of 4. Owing to scaling from model to model and to radical temperature corrections, this spacing is not always maintained. To compress the tables, all the variables except temperature are given as logs. The units are cgs for all variables. The columns are mass per unit area, Rosseland optical depth, continuum optical depth at 500 nm, geometric height, temperature, pressure, electron number density, atom number density, mass density, Rosseland mass absorption coefficient, continuum mass absorption coefficient at 500 nm, radiation pressure, radiative acceleration, and the fraction of flux carried by convection. The first values of $\tau_{\text{Ross}}, \tau_{\text{500}},$ and x are defined to be 0. Depths with no convection are listed with $\vec{0}$ convection fraction.

Kurucz 1979

Models depend on mass

- 100 M_{sun} < M: radiation pressure \rightarrow instabilities
- 0.08 < M < 100: hydrogen burning in the centre: **Main Sequence** $pc = 2...10^3$ g/cm³ <p> ~ 1
- \cdot 0.001 $<$ M $<$ 0.08: e gas degenerate in centre: no hydrogen fusion (low T): **brown dwarf**

 $pc = 10...10^3$ $$\rho$$

• M < 0.001: e– degenerate, solid core, **Jupiter**like $\rho_c \sim 10$ $\langle \rho \rangle = 1$

For comparison

- White dwarf $\epsilon_{\rm 0}$ $\epsilon_{\rm 0}$ \sim 10^6
-
- Black hole ϵ_0 \ge \sim 10¹⁷
- **Neutron star** $\leq \rho$ **>** $\sim 10^{13}$ **...10¹⁵**

MS stars: conditions at centre

Main sequence stars

Convective core = always well mixed Radiative envelope

CNO cycle: strongly conc. to centre

M > 20 Msun radiation.pressure

Radiative core

Convective envelope M<0.3 fully convective

pp chains: low conc. to centre

radiation pressure unimportant

Main sequence (solar compos.)

Lifetime on the Main Sequence

•
$$
\tau_{MS} = \frac{M_H}{|M_H|} \propto \frac{M_*}{L_*} \propto \frac{M_*}{M_*^3} \propto M_*^{-2}
$$

more accurate (from Geneva models)

 $\tau_{MS} = \frac{1}{10 \text{ Gyr}} \frac{(3 + M^{-1.6}) \text{ Myr}}{10 \text{ Gyr}}$ for M 10 Gyr /M³ for M 10 Msun

- Massive stars leave the main sequence earlier than less massive ones. Evolution after burnout is fast \rightarrow the turn-off mass is a unique function of the age of a stellar population
- Measures ages in absolute terms (nuclear rates)

Table of main sequence

Table of main sequence (II)

CMD open cluster: Praesepe

Haffner 1937

CMD globular cluster

CMD of clusters of different ages

After Eggen & Sandage 1962

Situation at end of Main Sequence

- Hydrogen exhausted in the centre
- Hydrogen burns in a shell around He-core which eats itself outward
- Gravity \rightarrow contraction of core + expansion of envelope (found by numerical sims.)

Solar evolution

Maeder 1989

Contraction of the core

Consider homologous contraction of gas sphere $(R,p,\rho,T$ with M=const):

- $\rho \propto MR^{-3}$ $\rightarrow \dot{\rho}/\rho = -3\dot{R}/R$
- \overline{dp} $\frac{d}{dt}$ = − $Gm\rho$ $\frac{dm\rho}{r^2}$ and $\frac{dm}{dr}$ $\frac{dr}{ }$ $= 4\pi r^2 \rho$ give $\frac{dp}{dr}$ \overline{dm} = − Gm $4\pi r^4$ so \overline{p} \overline{M} ≈ \overline{M} R^4 $\Rightarrow p/p = -4 \dot{R}/R$
- general EOS $\rho \propto p^{\alpha} T^{-\beta} \rightarrow \frac{\dot{\rho}}{2}$ ρ $=$ α \dot{p} \overline{p} $-\beta$ \dot{T} \overline{T}

$$
\bullet \quad \Rightarrow \quad \frac{\dot{T}}{T} = -\frac{4\alpha - 3}{\beta} \frac{\dot{R}}{R}
$$

Contraction of the core …

• Perfect gas $(\alpha = \beta = 1)$:

 \dot{T} \overline{T} = − \dot{R} \overline{R} contraction > core heats up

• Completely degenerate gas (non rel.) $\alpha = 3/5$ 0 < β < 1:

 \dot{T} \overline{T} $= +$ 3 5β \dot{R} \overline{R} contraction > core cools down

{This condition is already reached at some value of $v \sim \ln E_f/kT$, long before complete degeneracy!}

The H,He-burning processes

Fusion processes beyond He

Situation in the core

The smaller its mass, the closer a star is to degeneracy, the fewer burning phases it has:

M < 0.08 no H burning $M < 2$ no He burn. M < 9 no CO burn.

Maeder 1989

Massive stars (M > 10 Msun)

- Do all burning phases $(H, He, ... Si \rightarrow Fe)$
- Onion-shell structure + shell sources

Shell sources in M=25 Msun

Massive stars: final fate

After end of $Si \rightarrow Fe$ burning:

- Collapse of Fe core
- Collapse of envelope \rightarrow kinetic energy 10⁵¹ erg

0.3 s
$$
\left\{\n \begin{array}{l}\n \rho \sim 10^{11} \\
 \tau F F \sim 0.04 \text{ s}\n \end{array}\n\right.
$$

- Bouncing of envelope at centre: shock wave thru envelope \rightarrow compresssion \rightarrow ignition of fusion reactions
- = supernova explosion
	- Rapidly expanding envelope
	- Explosive nucleosynthesis in envelope
	- Collapse of core: neutron star/black hole

Explosive fusion

Hayashi line/tracks

- Locus in the HR diagram of all completely convective, hydrostatic stars of the same mass
	- and composition
- The low-T limit of hydrostatic stars: Any star cooler will evolve within ~100 d towards $HL \rightarrow$ stars can exist below HL for transitory phases only
- Examples
	- Protostellar collapse
	- HL is limit for giants

Hayashi 1961

Protostellar evolution tracks

Iben 1965

Evolution after Main Sequence

HRD and CMD solar nh'd

HRD evolutionary tracks

Maeder 1989

HRD – CMD – HRD

Evol.tracks & isochrones

Evolution: massive stars

- rapid evolution to the 'right' in the HRD
- Further burning phases are short
	- He: ~10 … 20% of MS-lifetime

 $-C:$ ~100 yrs

- Strong stellar winds (driven by radiation pressure of hot+luminous stellar atmosph.) modify evolution strongly: $|\dot{M}| \approx 10^{-6}$ Msun/yr τ ms ~ 3..10 Myr → star loses 3 ... 10 Msun
- Rotation

Evolution: massive stars

Intermediate mass stars (3..8 Msun)

- Core H exhaustion: core contracts & envelope expands
- Shell H burning
	- convection zone comes down (1st dredge-up)
	- move to Hayashi line = RGB
- Ascend on RGB
	- Mass loss (high $L \rightarrow$ radiatively driven wind)
	- At tip: core He-burning ignites

Evolution of 7 Msun star

Intermediate mass stars (3..8 Msun)

- Core expands & envelope contracts: **Blue Loop**
- back to Hayashi line = AGB
- Ascend on AGB
	- convection zone comes down (2nd dredge-up): brings down fresh hydrogen \rightarrow reignition of Hburning shell
	- mass loss
	- thermal pulses (TP) = instability of He-shell (shell flashes)

Thermal pulses

after Pols et al. 2001 after Pols et al. 2001

Thermal pulses (duration 10..100 yr, period 10³..10⁵ yr) 9 departure from AGB
¢= 0.89 8 sur 'CNO $\overline{7}$ Ήe 0.89 6 $\begin{array}{ccc}\n\text{Log}(L/L_\text{sun})\\
\omega & \rightarrow & \text{on}\n\end{array}$ 2 0 668.6 668.8 669.0 669.2 669.4 669.6 669.8 669.90 669.84 Time $[10^6]$ yr]

Pols et al. 2001 Pols et al. 2001

IMS evolutionary tracks

Ignition of central He-burning

Blue loops

IMS evolution

- Above ~5 Msun: Hot bottom burning: hot enough for CNO cycle \rightarrow ¹²C \rightarrow ¹⁴N
- Final fate:
	- $-$ Short strong wind or expulsion of outer envelope \rightarrow planetary nebula
	- Remnant star (CO core + 10-3 Msun H envelope) contracts at constant luminosity (core massluminosity relation) $\rightarrow T \sim 10^5$ K and ionizes nebula
	- Shell sources extinguish after ~30000 yr
	- Remnant star (= white dwarf) cools out (τ ~10 Gyr)

Low mass stars (0.08..3 Msun)

- After core H exhaustion to Hayashi Line = RGB
- Ascend on RGB: at the tip if M_{He} > 0.45 M_{sun} ignition of He:
	- Degen.matter cannot compensate heat input by expansion \rightarrow T increases \rightarrow burning increases: thermonuclear runaway = **core He-flash** (timescale ~hours)
	- T increases until degeneracy is overcome: He burns steadily
	- He-burning shell eats down into core (~10⁵ yrs) \rightarrow core He-burning (**Horizontal branch**)

Horizontal Branch stars

• Position depends on chemical composition and on 2nd parameter (age, core mass …)

• Exhaustion of core He supply \rightarrow AGB (but without 2nd and 3rd dredge-ups)

Problems & difficulties

- Opacities:
	- need absorption cross sections from all levels of all ions of all elements (OPACITY Project)

– Summation over all relevant levels

- Energy production rates: 1982 rate for ${}^{12}C(\alpha,\gamma)$ ¹⁴N increased by factor 3..5
- Equation of state: difficult at high T and ρ (missing energy levels, atoms change structure due to vicinity of other atoms)

• convection

No physical theory available

Use simple approaches (mixing length 'theory') and adjust their fudge parameters (mixing length l/Hp)

Use formulae derived from numerical simulations of convection (now possible in 3D)

Problems & difficulties

- Stellar winds (recipes from observations or theory)
- Dynamical phases (He-flash, TP, PN, **SN**)
- Double stars (mass transfer, SN Ia)

Rotation

- 3D models necessary; now possible
- **Large scale mixing** effects (Meridional circulation, various instabilities, ...)

Meynet+Maeder 2002

- Consequences of internal mixing
	- Slightly longer MS life times (20…30%)
	- Nucleosynthesis ('primary N' in very fast rotators … cf. later!)

Zero-metallicity stars (**Pop. III**)

