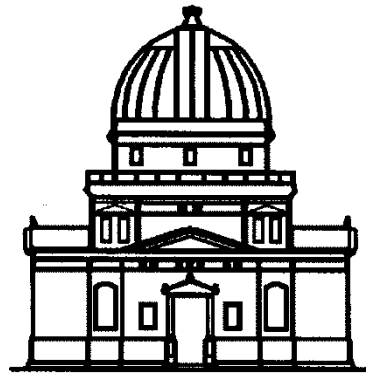


# Evolution of Galaxies:

## Galaxy formation, monolithic evolution



Observatoire astronomique  
de Strasbourg

J.Köppen      joachim.koppen@astro.unistra.fr

<http://astro.u-strasbg.fr/~koppen/JKHome.html>

# Galaxy formation

= instability of self-gravitating protogalactic cloud  
... like star formation

Dynamics of (isothermal) spherical cloud

$$\frac{\partial \rho}{\partial t} + \nabla(\rho \vec{v}) = 0$$

conserv. of mass

$$\frac{\partial \rho \vec{v}}{\partial t} + \nabla(\rho \vec{v} \cdot \vec{v}) = -\nabla p - \rho \nabla \Phi$$

conserv. of momentum

$$p = p(\rho, T) \propto \rho c^2$$

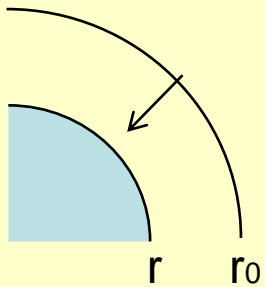
equation of state

$$\nabla \Phi = 4\pi \rho$$

gravitation

**Collapse if**  $M < M_J \propto T^{3/2} \rho^{-1/2}$  Jeans criterion

Free-fall solution (neglect p, OK for initial phase)



Mass shell initially at  $r_0$

$$t_{\text{ff}} = \sqrt{\frac{3\pi}{32G\rho}}$$

$$\frac{\pi t}{2 t_{\text{ff}}} = \sqrt{\frac{r}{r_0} \left(1 - \frac{r}{r_0}\right)} + \arcsin \sqrt{1 - r/r_0}$$

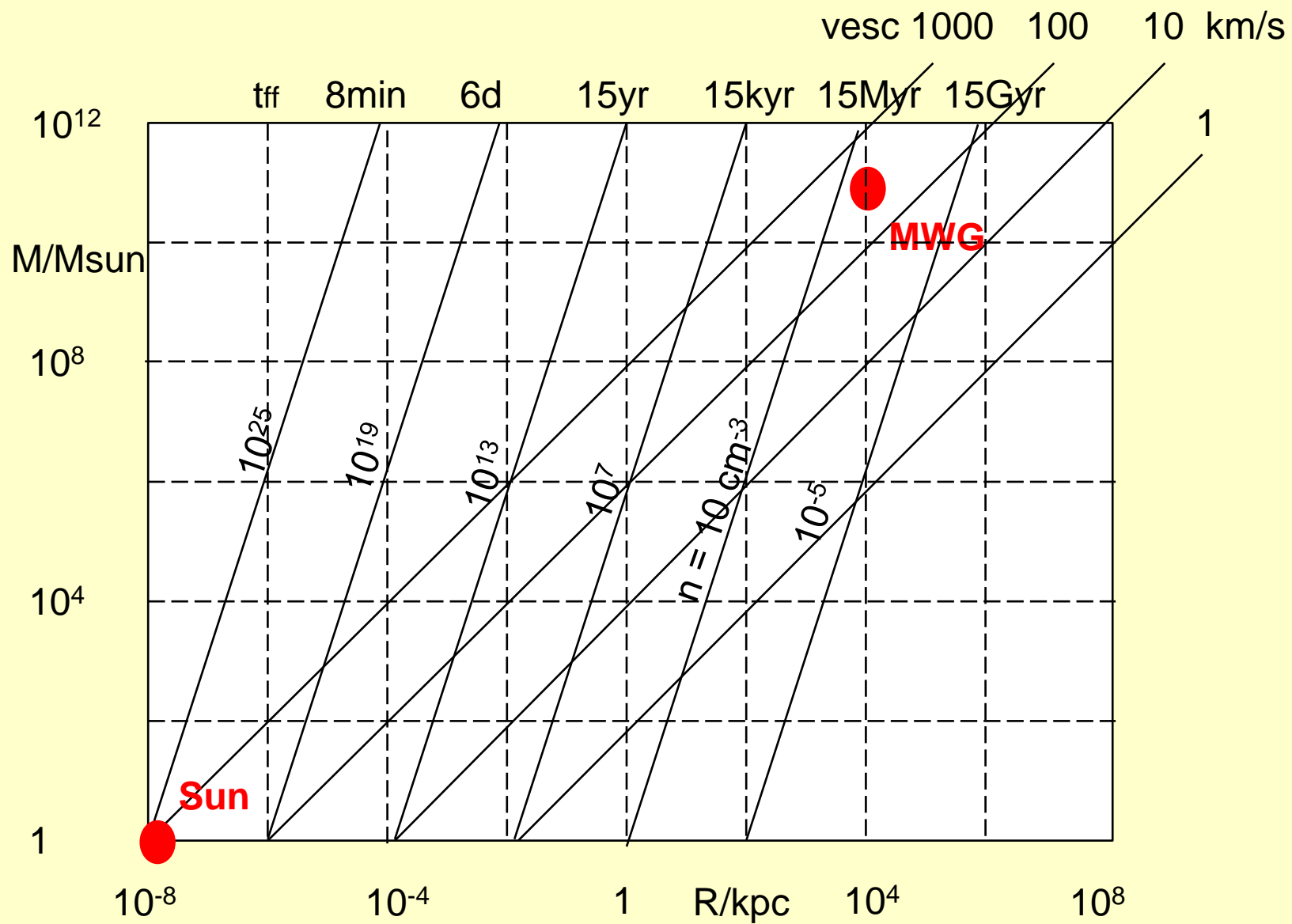
for  $t \ll t_{\text{ff}}$

$$r(t) = r_0 \left(1 - \left(\frac{\pi t}{4 t_{\text{ff}}}\right)^2\right)$$

$$v(t) = -\frac{\pi^2}{8} r_0 \frac{t}{t_{\text{ff}}} \propto -r_0 t$$

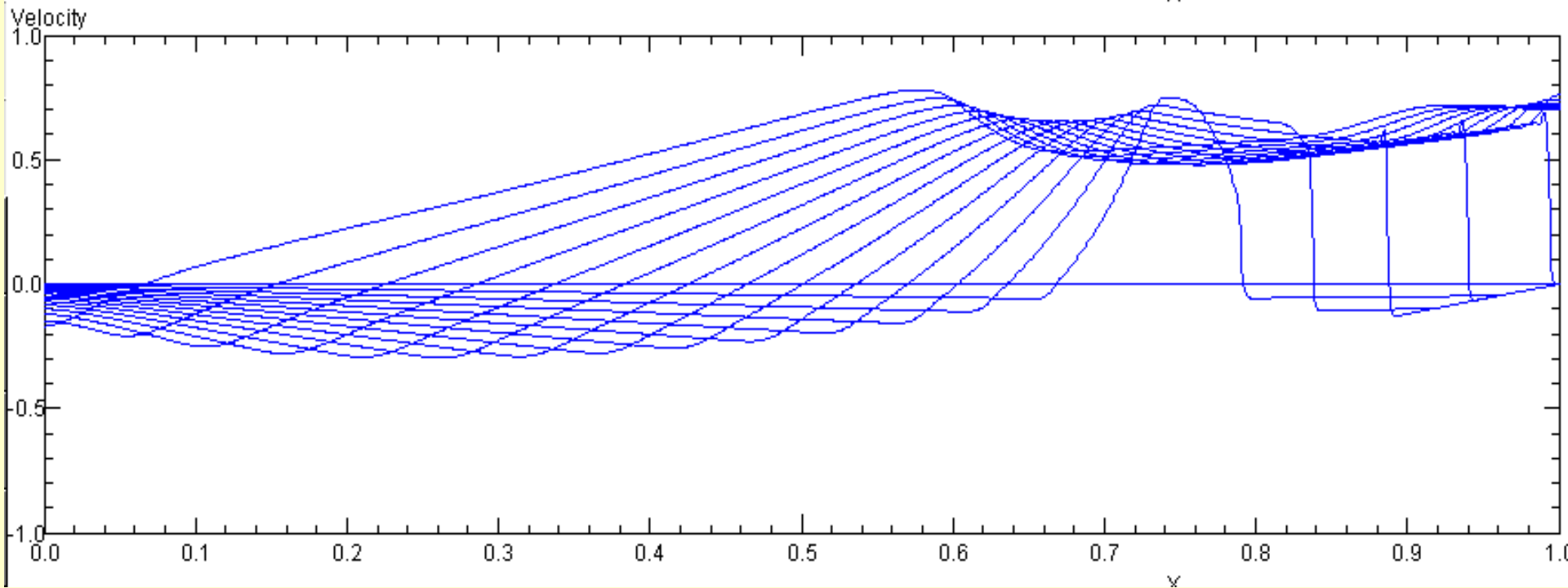
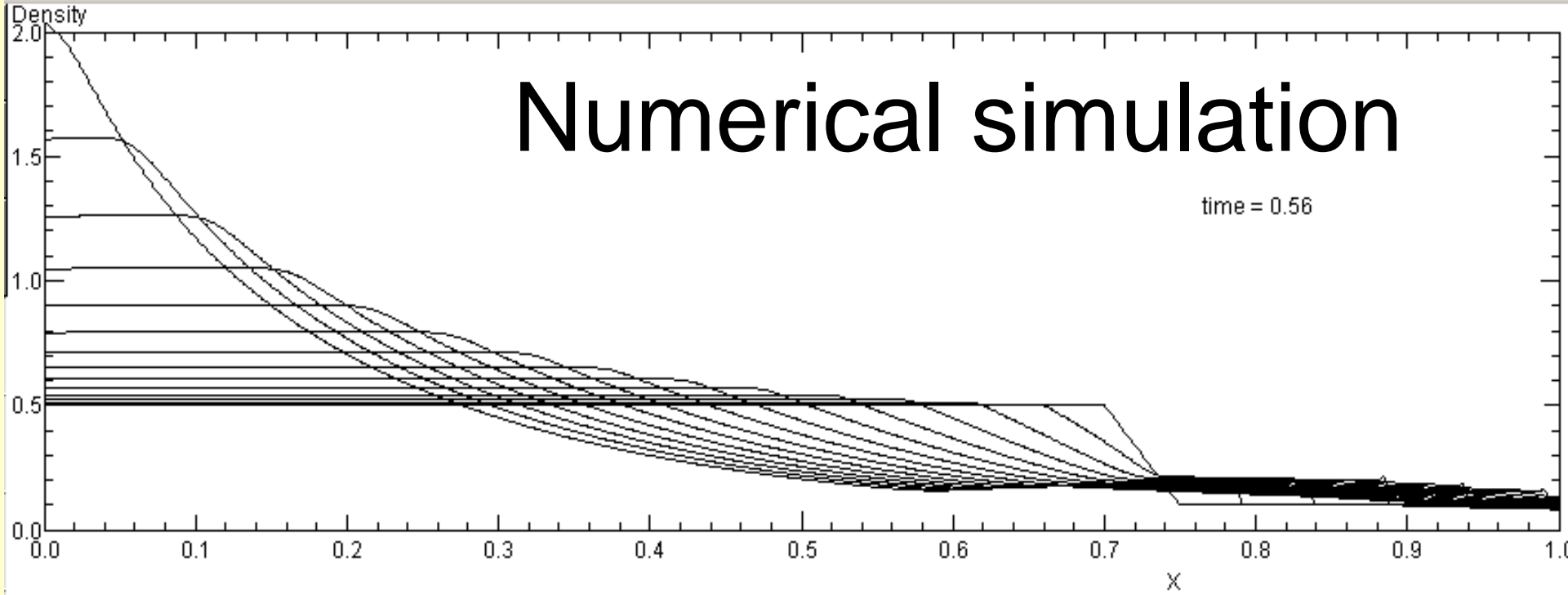
Complete collapse

$$r(t) = r_0 \left(1 - \frac{t}{t_{\text{ff}}}\right)^2 \rightarrow 0 \quad \text{for } t \rightarrow t_{\text{ff}}$$



# Numerical simulation

time = 0.56

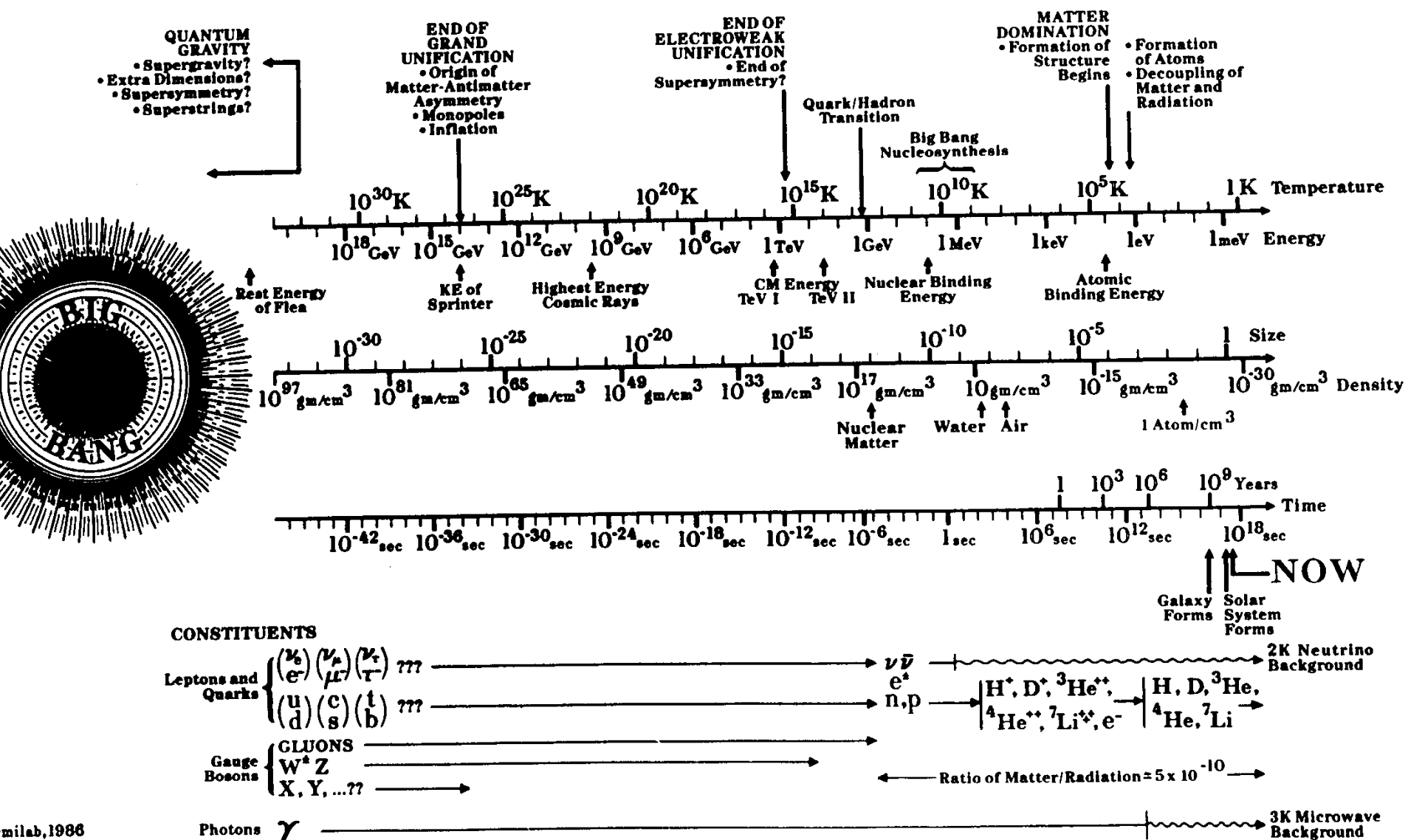


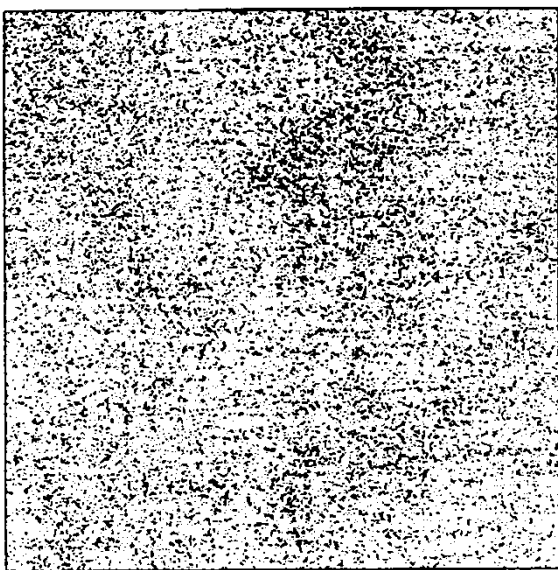
# Galaxy formation in the Universe

- Big Bang

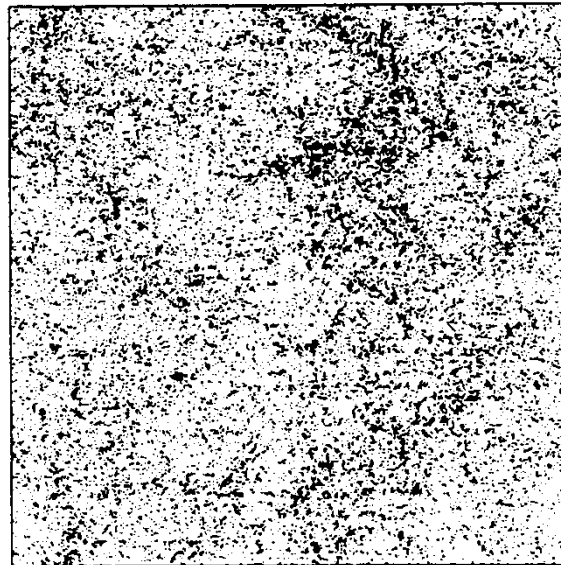
- at  $T = 10^5$  yrs matter/radiation decoupling
- Hot gas cools by expansion and radiative cooling
- Density fluctuations grow if self-gravitation overcomes expansion
  - Condensation can collapse if  $M > M_J$
  - cools and forms stars, if cooling time  $<$  free fall
- Cooling curve  $\rightarrow T \sim 10^4$  K  $\rightarrow M_J \sim 10^7 \dots 10^9$  Msun  $\rightarrow$  size of fragments

# History of the Universe

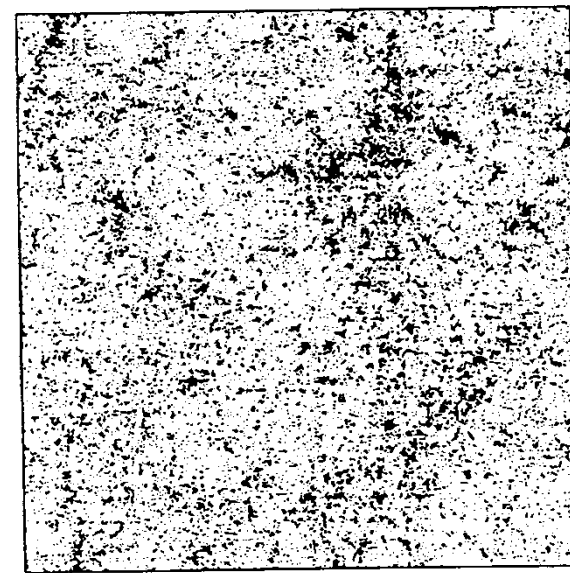




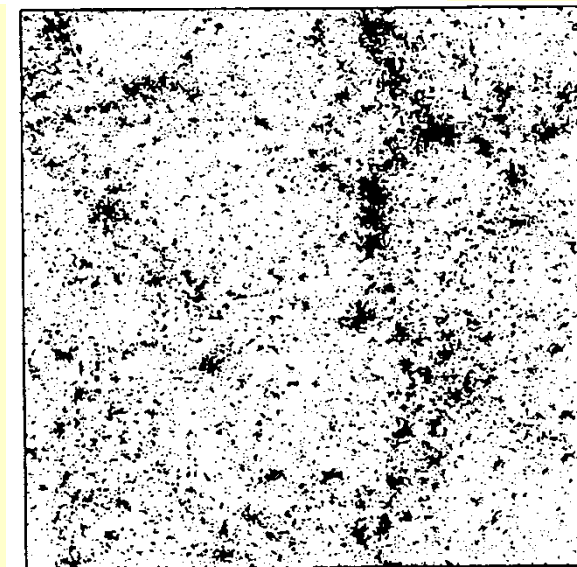
$R = 1$



$R = 1.8$



$R = 2.4$



$R = 4.5$

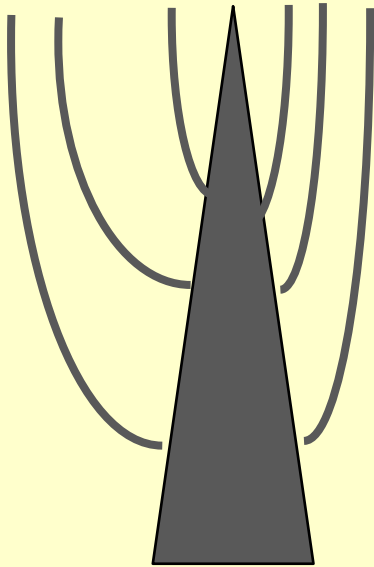
Formation of (dark matter) structures  
By the gravitational instability in  
an expanding universe:

Denser region attracts mass,  
becomes denser and attracts  
even more mass ...

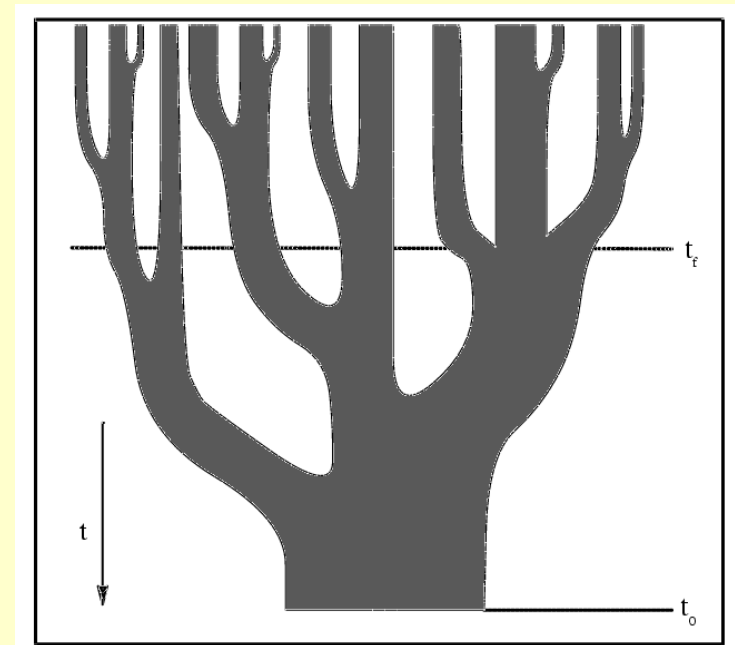


# Fate of Fragments

- Fragments are building blocks (~ dwarf galaxies) from which galaxies assemble by merging



Spiral (always prominent core)



Elliptical (equal mass partners)

# Dissipative (Monolithic) Collapse

- Larson (1969ff): protogalaxy consists of Jeans-unstable clouds ('turbulence elements') whose population can be treated by the Boltzmann equation (moments  $\rightarrow$  fluid equations)
- Let's consider a spherical protogalactic cloud ...

# For the gas clouds ...

$$e = \frac{1}{2}\rho(\sigma_r^2 + \sigma_\varphi^2 + \sigma_\vartheta^2) = \frac{3}{2}\rho\sigma^2 = \frac{3}{2}p \quad \text{Gas is isotropic}$$

$$\frac{\partial \rho}{\partial t} + \nabla(\rho \vec{v}) = -\Psi$$

$$\frac{\partial \rho \vec{v}}{\partial t} + \nabla(\rho \vec{v} \cdot \vec{v}) = -\vec{v}\Psi - \nabla p - \rho \nabla \Phi$$

$$\frac{\partial e}{\partial t} + \nabla(e \vec{v}) = -\frac{e}{\rho}\Psi - p \nabla \vec{v} \quad - (de/dt)_{\text{diss}}$$

# Stars (collisionless Boltzmann →)

$$\sigma_{\text{rad}}^2 \neq \sigma_{\text{tang}}^2 = \frac{1}{2}(\sigma_{\varphi}^2 + \sigma_{\vartheta}^2) \quad \text{can be anisotropic}$$

$$\frac{\partial \rho}{\partial t} + \nabla(\rho \vec{v}) = \Psi$$

$$\frac{\partial \rho \vec{v}}{\partial t} + \nabla(\rho \vec{v} \cdot \vec{v}) = \vec{v}_g \Psi - \rho \nabla \Phi - \frac{\partial \rho \sigma_r^2}{\partial r} - \frac{2\rho}{r}(\sigma_r^2 - \sigma_t^2)$$

$$\frac{\partial \rho \sigma_r^2}{\partial t} + \nabla(\rho \sigma_r^2 \vec{v}) = -\left(\frac{e}{\rho}\right)_g \Psi - \frac{2\rho \sigma_r^2}{r} \frac{\partial v_r}{\partial r}$$

$$\frac{\partial \rho \sigma_t^2}{\partial t} + \nabla(\rho \sigma_t^2 \vec{v}) = -\left(\frac{e}{\rho}\right)_g \Psi - \frac{2\rho \sigma_t^2}{r} \frac{v_r}{r}$$

Advection

interaction

---

Source terms

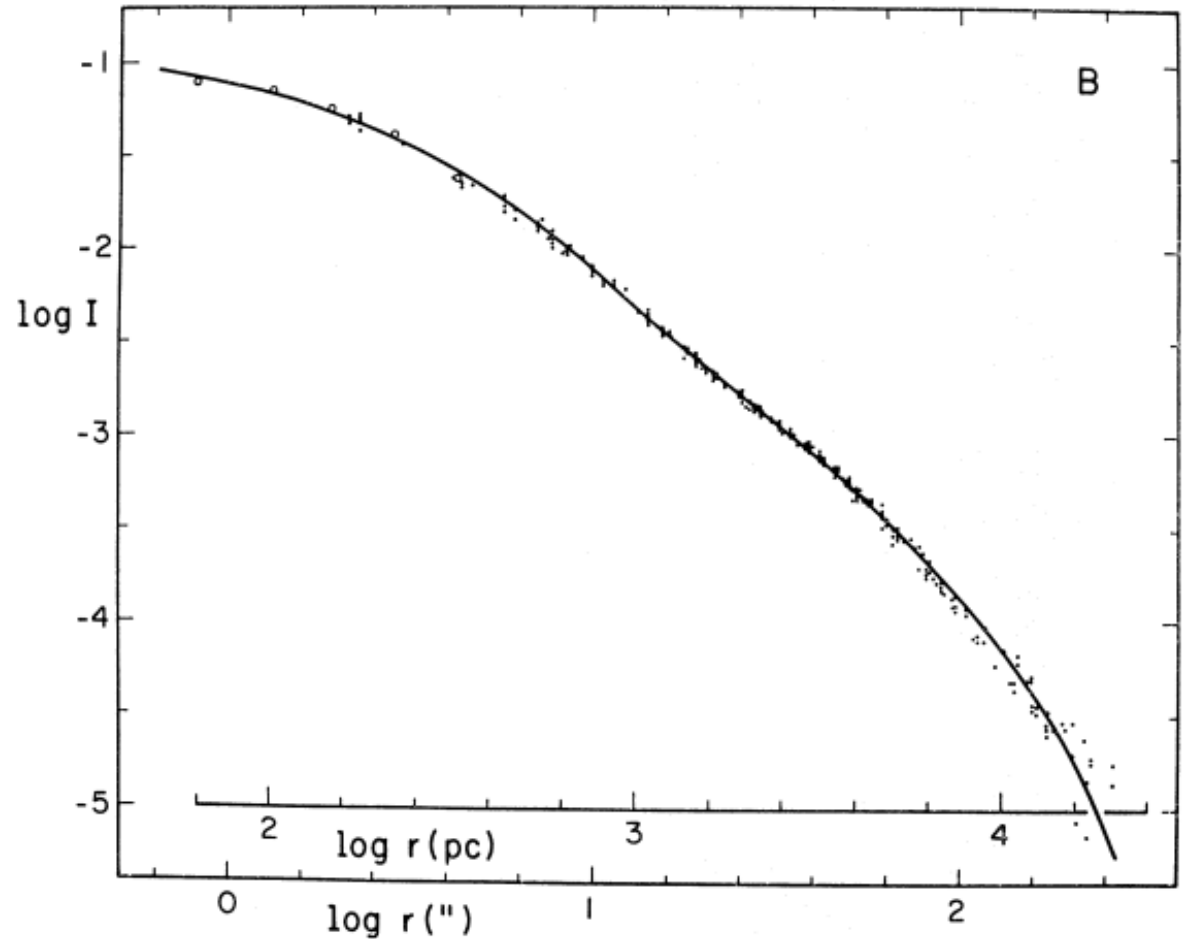
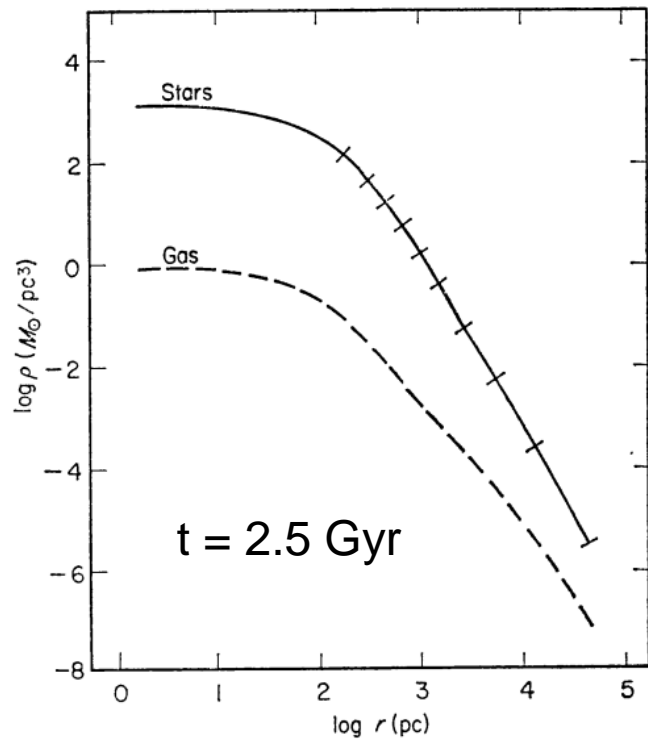
# Assumptions (Larson 1969)

- $M = 10^{11} \text{ Msun}$ ,  $R=50 \text{ kpc}$ ,  
 $\rho = 2 \cdot 10^{-4} \text{ Msun/pc}^3 \rightarrow t_{\text{ff}} = 600 \text{ Myr}$
- $de/dt \propto \rho^{1/2}$  collisions of clouds
- $\text{SFR } \Psi \propto \rho^n$   $n=2$  cloud-cloud collisions  
(reasonable guess)

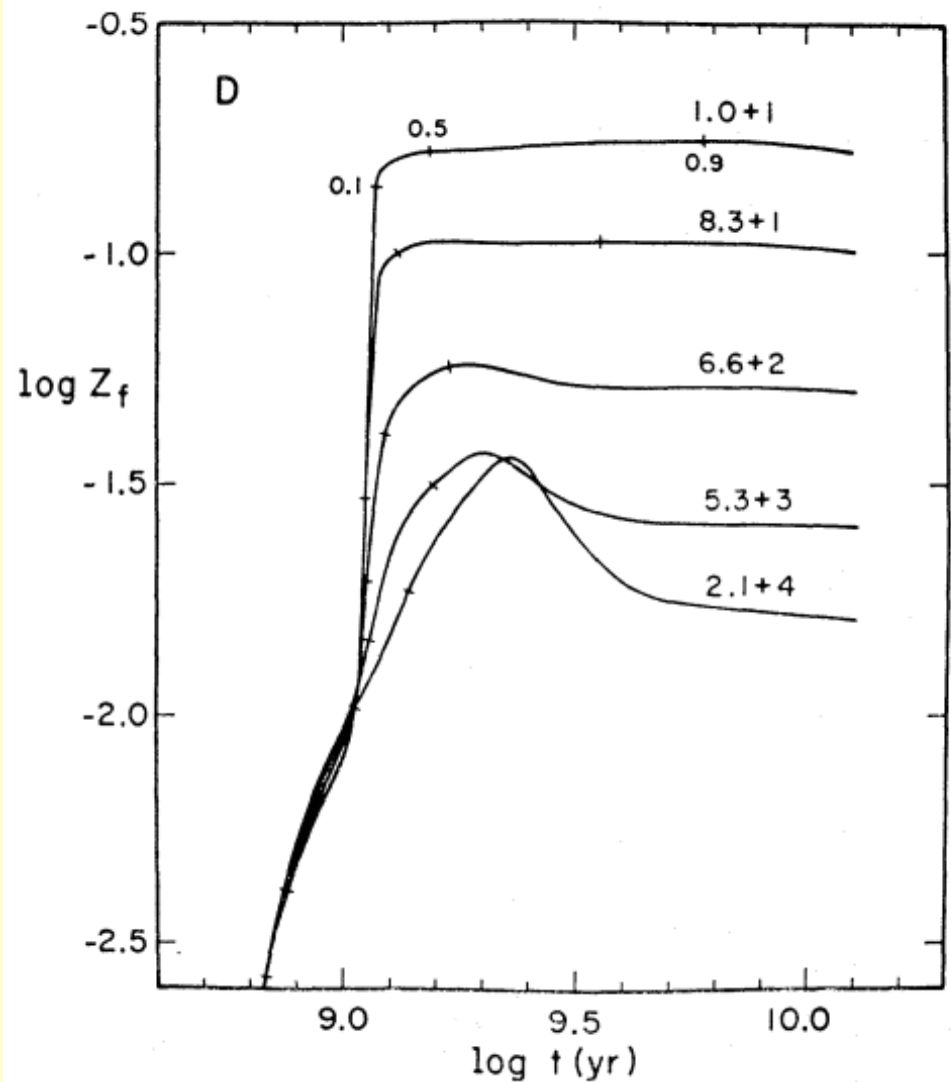
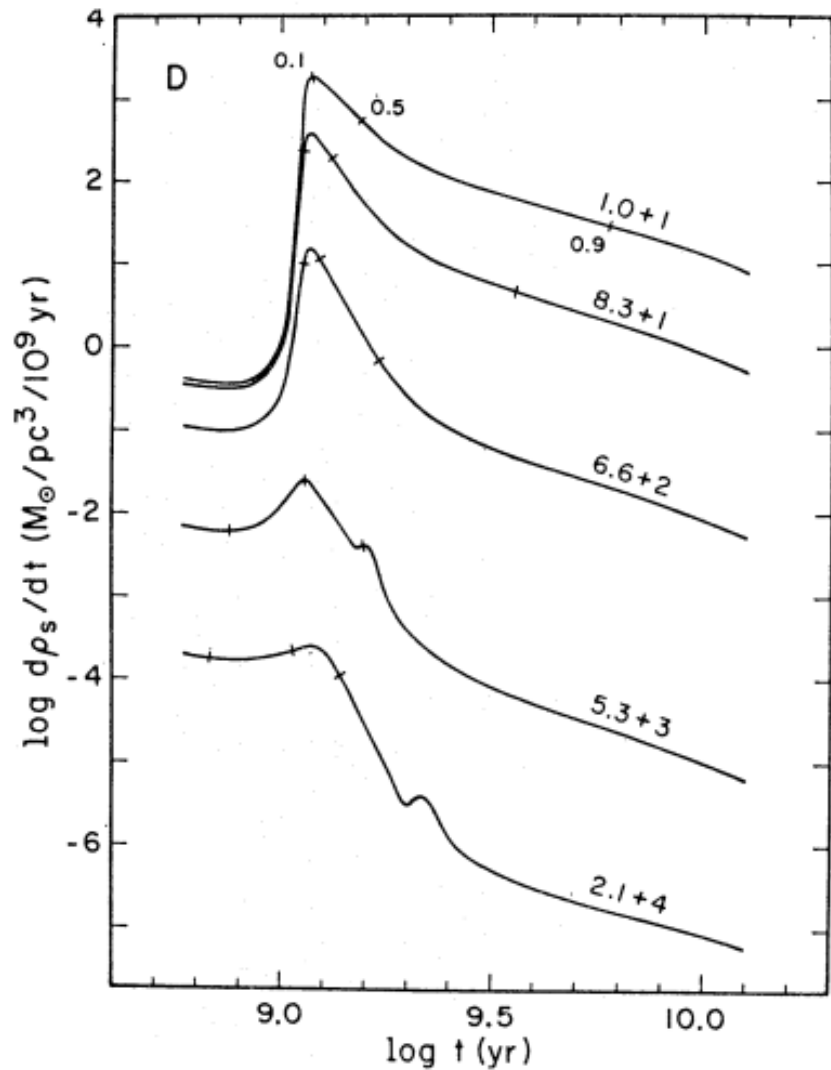
# Results (Larson '69 + '74)

- 800 Myrs:  $\rho_{\text{gas}} = \rho_{\text{stars}}$  at centre (1 atom/cm<sup>3</sup>)
- 1.7 Gyrs: 50 at/cm<sup>3</sup>; 97% stars
- 2.5 Gyrs: 99.6% stars, stopped (numerics)
  
- Stellar density profile  $\rho \propto r^{-3}$  OK!
- Stellar velocities nearly isotropic
- Metallicity gradient: negative
- Collapse only if SFR power  $n > 1.8$

# Stellar profile OK with observation!

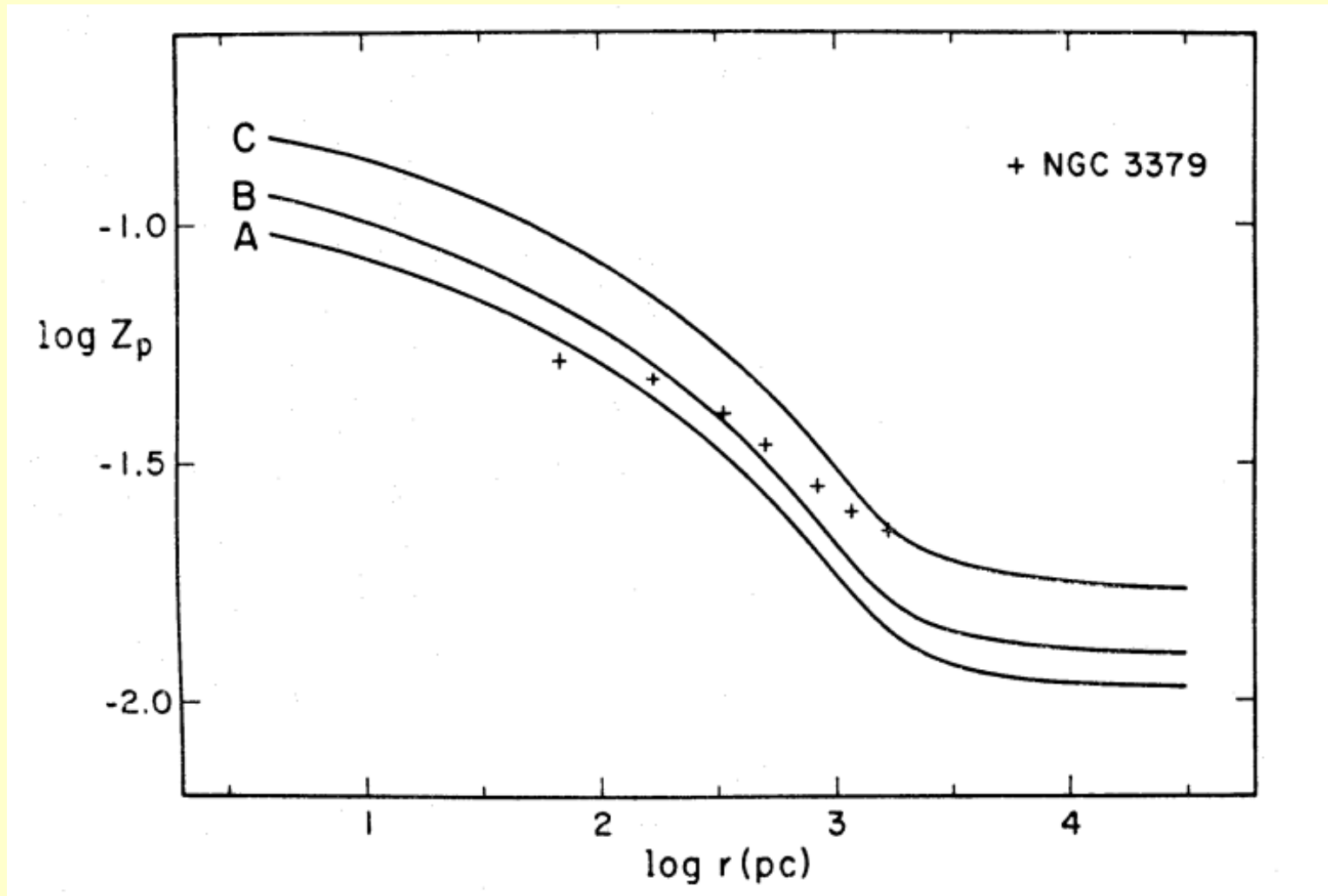


# SFR and metallicity (1974)





# Projected metallicity profile



# However:

- The centre has higher proportion of younger stars → the centre is bluer by about 0.2 mag
- Observations show that centre is redder!

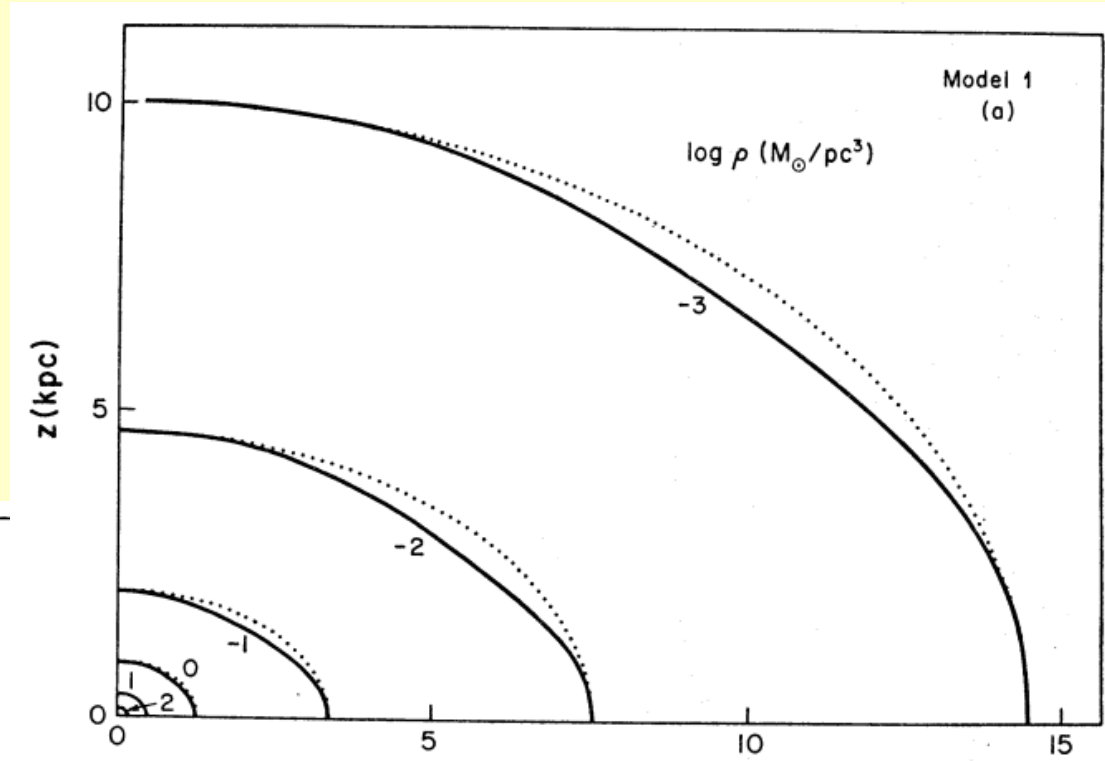
## Remedy:

- Expell gas by SN driven terminal wind → cuts star formation → can account for observed photometry and the mass-colour relation

# Rotating Ellipticals (1975)

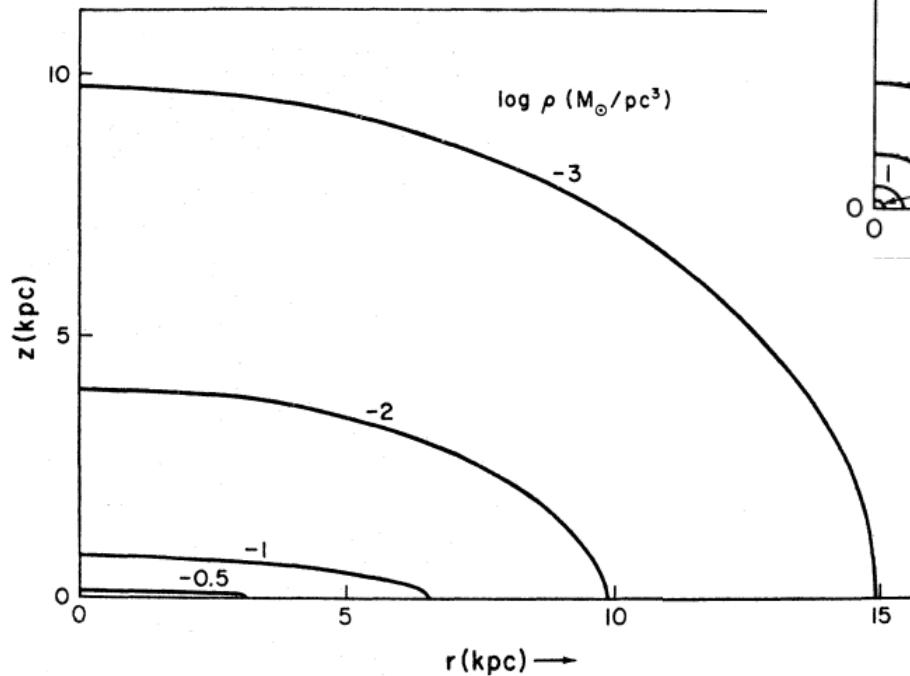
- Not all E are E0
- 2D models
  - Formation of disk in equatorial plane → isophotes would NOT have same ellipticity as it is observed
  - Remedy: postulation of strong viscosity

# Rotating Ellipticals (1975)

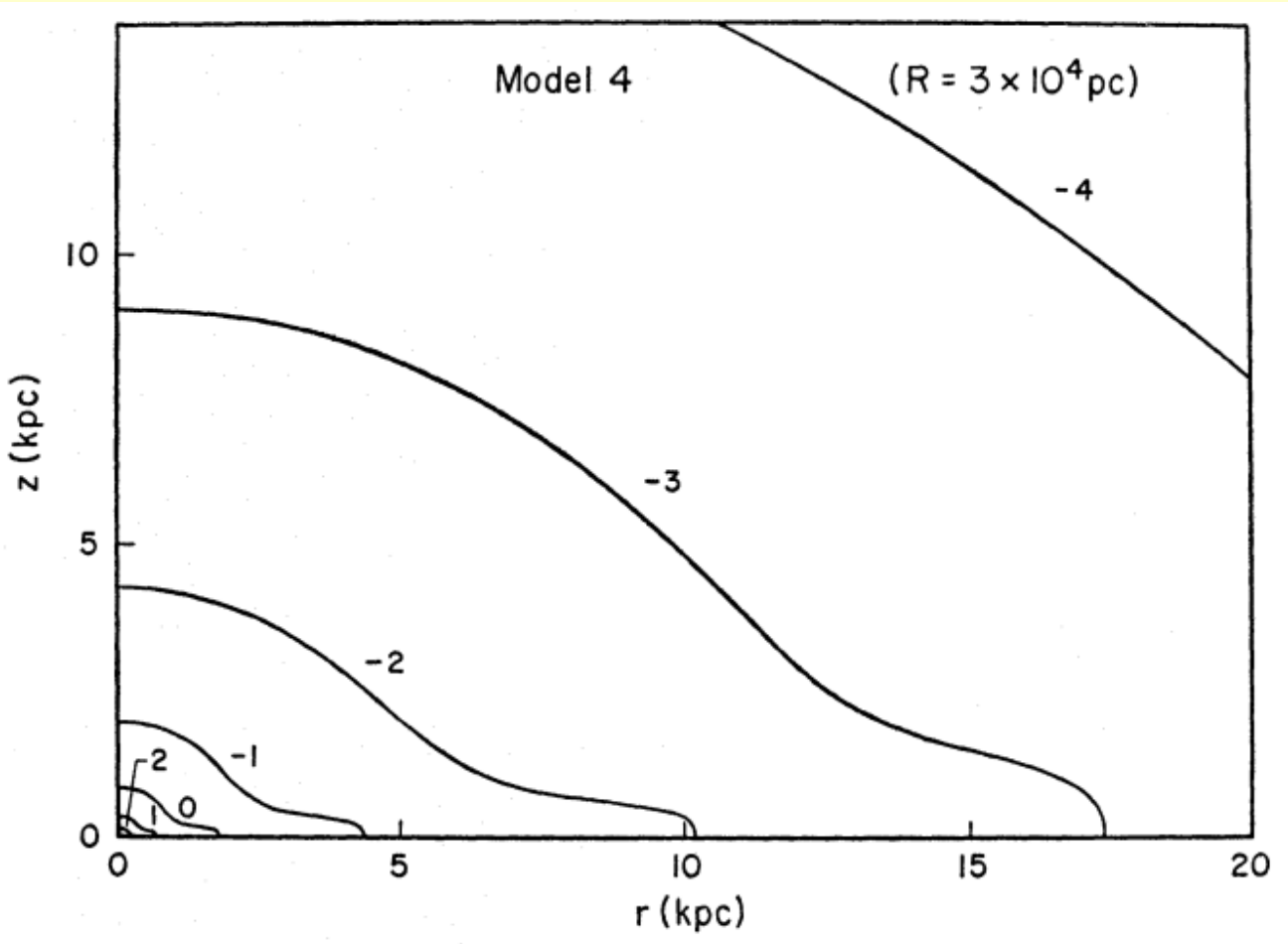


With extra viscosity

Density contours: no viscosity



# Disk galaxies (1976)



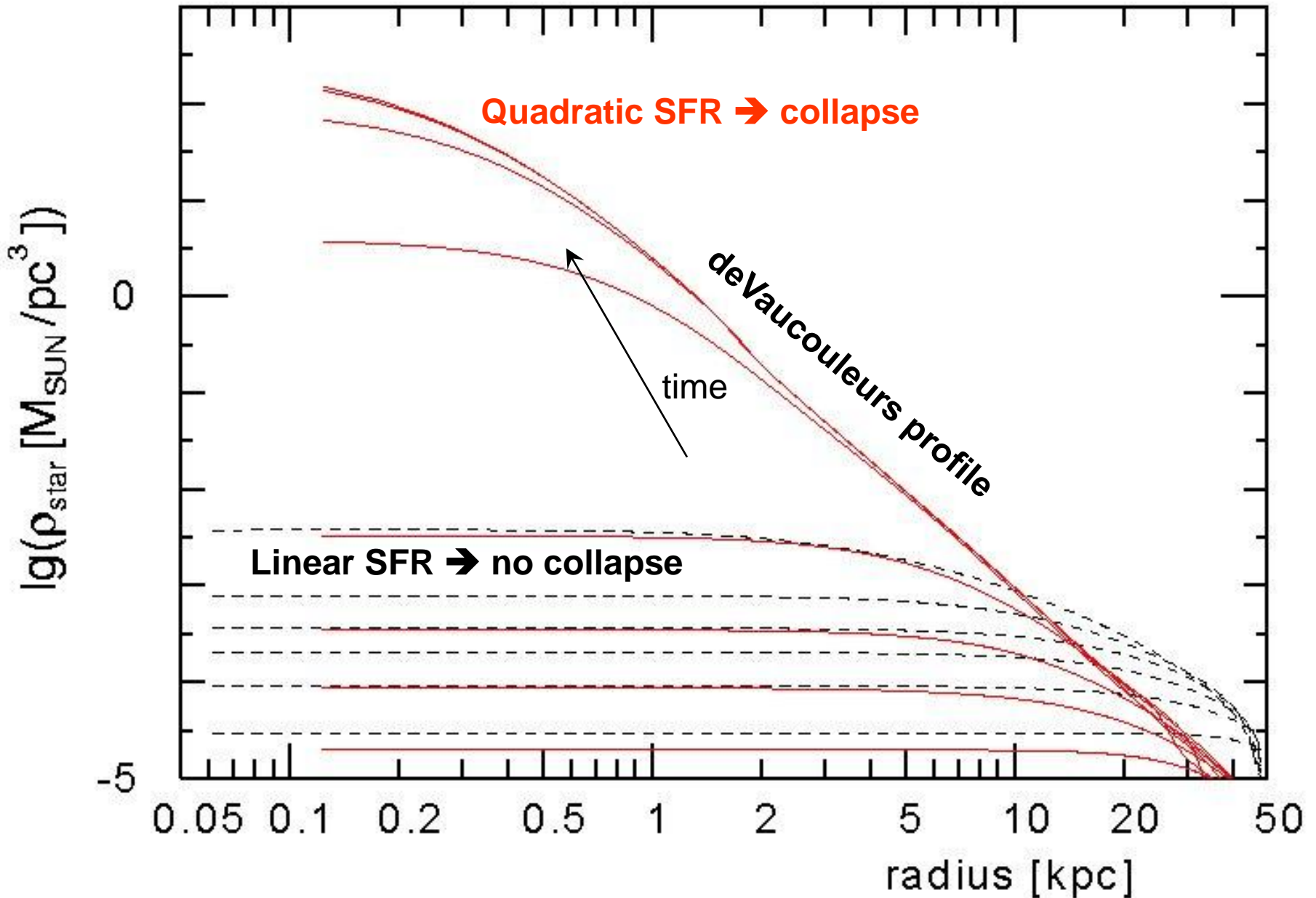
Needs time-dependent  
viscosity and SFR



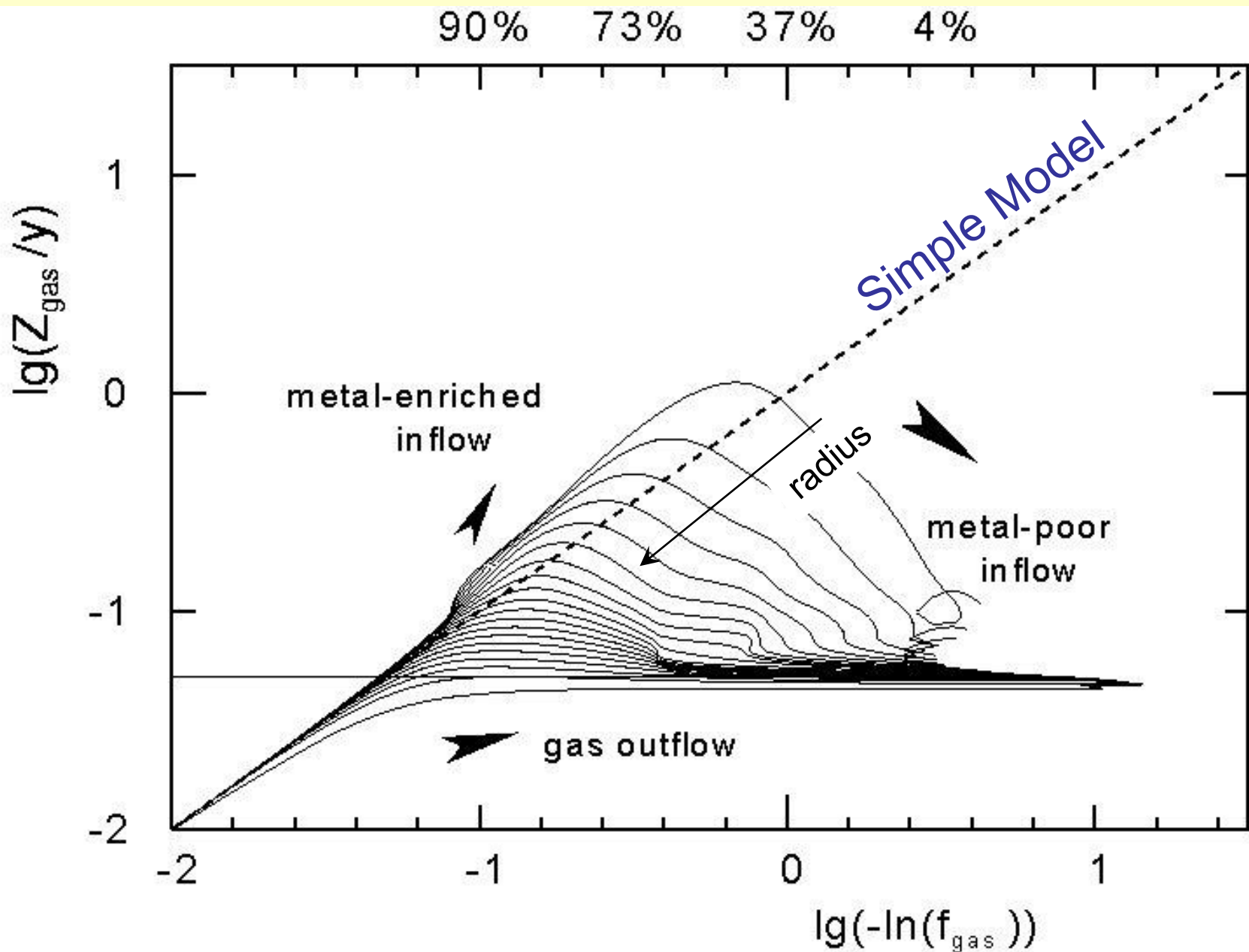
# An alternative

- Dissipation-less (stellar dynamical) models of ellipticals:
  - N-body simulations for stellar population (Gunn 1976)
  - Nice density profile is obtained, if a heavy halo ( $6 \cdot M_{\text{gal}}$ ) is assumed
  - does not explain metallicity gradient ...

# Another look at spherical collapse



# Another look at spherical collapse





# The G dwarfs ...

